RAMAN SPECTROSCOPY OF VIBRATIONAL AND ROTATIONAL RELAXATION OF ACETONITRILE MOLECULES DISSOLVED IN IONIC LIQUIDS

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Vibrational and rotational relaxation processes of acetonitrile (AN) molecules as molecular probes dissolved in the imidazolium-based room-temperature ionic liquid (RTIL) 1-decyl-3-methyl-imidazolium bromide have been studied by Raman spectroscopy technique (the band shape analysis using different polarizations of the incident and scattered light). It has been shown that the vibrational relaxation processes are quickened in the RTIL surrounding. The corresponding vibrational correlation times decrease from about 3 ps in the neat AN to 1.5–2 ps in AN solution in RTIL (10 and 20 wt % of AN). The effect of RTIL environment on the rotational dynamics of AN, namely, on the tumbling motion of dipole moments, is rather small. However, the slow-down of AN reorientations in RTIL was distinguished, namely, the rotational correlation times increase from 1.0–1.5 ps (neat AN) to 1.6–1.9 ps (AN/RTIL solutions). The results have been discussed in terms of various molecular and media contributions to these processes (H-bonding, highly viscous media effect, intermode exchange of vibrational energy, etc). The demolition of structural effects (short-range order and mesoscopic structures in the nano-scale) in 1-decyl-3-methyl-imidazolium bromide by the mutual penetration of AN and RTIL species is suspected.

Keywords: properties of molecules and molecular ions, line and band widths, shapes and shifts, ionic liquids

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1. Introduction

Ionic liquids (ILs) are one of the red-hot and most successful breakthroughs in creating multifunctional materials. ILs have already demonstrated their unique features in many fields of high technology. Particularly notable are enzyme catalysis, protein synthesis, membrane technology, battery and fuel cells, dye-sensitized solar cells (DSSC), gas sensors, nanotechnology, etc (see Refs. [1–3] and the literature therein). Nevertheless in this wealth of work certain important gaps remain. Namely, investigations of the fundamental molecular processes, such as conformational transitions, phase separation, proton transfer, vibrational and rotational relaxations occurring in various IL environments. Note that due to their dielectric properties (static dielectric constants of various classes of ILs are spread usually over 11–14 [4, 5]) the IL systems are not ‘super-polar’. Their extraordinary action in certain chemical reactions and in other molecular processes is mainly due to supra-molecular structuring in nanoscale effects and possible equilibria within the IL phase [3]. Hence, understanding of the peculiar role of ILs as solvents is not possible without understanding their microscopic properties, e.g. short-range ordering, microscopic friction effects, etc. These can be studied by the micro-dynamics of molecular probes, i.e. measuring correlation times of vibrational and rotational relaxation of spherical (neopentane, tetramethylsilane, etc) or symmetrical top (acetonitrile, chloroform, etc) molecules dissolved in various IL systems. Since those processes run at ordinary conditions in the time scale of picoseconds, a Raman spectroscopy technique (the band shape analysis using different polarizations of the incident and scattered light, i.e. isotropic and
anisotropic band shape analysis [6, 7]) can be successfully exploited.

The purpose of the present work was investigation of vibrational and rotational relaxation processes of acetonitrile molecules as molecular probe dissolved in the room-temperature ionic liquid 1-decyl-3-methyl-imidazolium bromide by Raman spectroscopy.

2. Experiment

Commercially available 1-decyl-3-methyl-imidazolium bromide (abbreviation: aha2/cdi16, CAS Nr 188589-32-4, Fig. 1) of \( \geq 98\% \) purity from Merck KGaA Darmstadt and acetonitrile (AN) of high purity (99.5 \( +\) %) from Aldrich Chemical Co were used in the present work without additional purification. The samples were prepared by weighing (\( \pm 0.1 \) mg) the components.

Raman spectra were recorded using a DFS-12 double-grating spectrometer with a slit width of 0.08 mm. Similarly as in [8], this value of the slit width was found to be optimal. Namely, registration of the weakest Raman bands was still possible at this setting, whereas the instrumental broadening of the observed bandwidth could be neglected. The excitation source was an argon-ion laser (Stabilite, SpectraPhysics) operating at 488.0 nm (120 mW) or 514.5 nm (120 mW). All measurements were carried out at the 90° scattering geometry. The required polarization of the incident and scattered light was selected by a conventional polarizer. The detection part of the instrument consists of photo-multiplier and a photon-counting system online with a standard PC. The Raman spectra band shapes were processed digitally using the Microcal Origin [9] program package.

3. Results and discussion

We have chosen for the present study one of imidazolium-based room temperature ionic liquids (RTILs), namely – 1-decyl-3-methyl-imidazolium bromide. This substance has a rather low melting point (289 K) and is totally miscible with water and acetonitrile. Perhaps it is the most intriguing that just above its melting point 1-decyl-3-methyl-imidazolium bromide is extremely viscous (7892.9 mm²/s at 293 K, all data from Merck KGaA [1]). Thus an extreme frictional action and hence a peculiar dynamics of solute molecules may be expected in this RTIL as a solvent.

The Raman spectrum of 1-decyl-3-methyl-imidazolium bromide is shown in Fig. 2. Several free spectral ‘windows’ are easily noticeable, namely below about 600 cm\(^{-1}\), rather narrow spots at 900 and 1200 cm\(^{-1}\), and a very broad region between 1500 and 2900 cm\(^{-1}\) that can be used analysing the band shapes of properly chosen probe molecules.

The symmetrical top molecules of acetonitrile were chosen as the probes in the present work for several serious reasons. First of all the rotational and translational dynamics of acetonitrile molecules in the neat AN as well as in various solvents have been studied using a wide kit of experimental techniques (FTIR, Raman, NMR relaxations, etc, see [8, 10–12] and the references therein). These data are very useful for comparative and metrological purposes. Furthermore, Raman bands \( \nu 2 \) (2249 cm\(^{-1}\)) and \( \nu 4 \) (918 cm\(^{-1}\)) usually used in those studies fit into the free spectral ‘windows’ of 1-decyl-3-methyl-imidazolium bromide (Fig. 2). Also note that AN molecules were successfully used probing reorientational motion in polyacrylamide gels [13].

The Raman band shape analysis of these AN modes (\( \nu 2 \) and \( \nu 4 \), \( C_{3v} \) point group symmetry and both modes of \( A_{1} \) type) is convenient also from the theoretical point of view. Namely, starting from a quite general treatment [6, 7], the isotropic and anisotropic profiles of Raman spectra \( I_{\text{iso}}(\omega) \) and \( I_{\text{anis}}(\omega) \) are related to the experimental shapes in VV and VH scattering geometry,
Fig. 3. Raman spectra (VV and VH) of neat AN in the ranges of modes $\nu/2 (2249$ cm$^{-1}$) and $\nu/4 (918$ cm$^{-1}$).

$I_{VV}(\omega)$ and $I_{VH}(\omega)$ respectively, via their convolution with the instrumental function $A(\omega)$ as

$$I_{VV} - \frac{4}{3} I_{VH} = A \otimes I_{iso},$$

$$I_{VH} = A \otimes I_{anis}. \quad (1)$$

If the slit widths are set properly, and thus instrumental disturbances of the band shapes can be neglected, and if the contours appear during processing to be Lorentzians having half-widths at half heights (HWHH) $\Delta$, the corresponding vibrational and rotational correlation times are calculated simply:

$$\tau_{vib} = \frac{1}{2\pi c \Delta_{iso}} \quad \text{and} \quad \tau_{rot} = \frac{1}{2\pi c (\Delta_{anis} - \Delta_{iso})}. \quad (2)$$

The information on the rotational dynamics obtained analysing $\tau_{rot}$ in the case of AN (i.e. C$_{3v}$ point group symmetry and A$_1$ mode) allows of a quite obvious interpretation. Applying the symmetry considerations [7], the rotational correlation functions can be expressed via single stochastically averaged Wigner $D$-function:

$$C_{\text{rot}}(t) = \langle D^2_{00}(\delta\Omega(t)) \rangle = \langle 3 \cos^2(\beta(t)) - 1 \rangle, \quad (3)$$

where $\delta\Omega(t)$ is the set of Eulerian angles that the molecule-fixed coordinate system (MCS) reorients after $t$ s with respect to the laboratory-fixed coordinate system (LCS) and $\beta$ is the polar angle of the principal symmetry axis of the molecule with respect to the LCS. It means that $\tau_{rot}$ obtained from the band shapes of $\nu/2$ and $\nu/4$ modes characterizes the tumbling motion (only $\beta$ changes) of AN molecules.

The Raman spectra (VV and VH) of the neat AN in the regions of $\nu/2 (2249$ cm$^{-1}$) and $\nu/4 (918$ cm$^{-1}$) modes are shown in Fig. 3. Secondary band structures, noted and considered in some earlier studies [9, 11, 13], are observed at the present instrumental setting as well. Asymmetric shoulders are attributed in some works to the “hot bands” [13], whereas in others [8, 10] those extra bands are believed to originate from intermolecular effects (aggregation of AN or H-bonding with solvent molecules). Without getting deep into polemic at
Table 1. Measured half-widths (HWHH) of Raman modes $\nu_2$ and $\nu_4$ and calculated vibrational and rotational correlation times ($\tau_{vib}$ and $\tau_{rot}$) in neat AN at $T = 295$ K and in AN solution in 1-decyl-3-methyl-imidazolium bromide (10 and 20 wt% of AN). Available literature data are given for comparison.

<table>
<thead>
<tr>
<th></th>
<th>$\Delta_{VV}$, cm$^{-1}$</th>
<th>$\Delta_{VH}$, cm$^{-1}$</th>
<th>$\tau_{vib}$, ps</th>
<th>$\tau_{rot}$, ps</th>
</tr>
</thead>
<tbody>
<tr>
<td>Neat AN</td>
<td></td>
<td></td>
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<td></td>
</tr>
<tr>
<td>$\nu_2$</td>
<td>1.65±0.08</td>
<td>6.10±0.97</td>
<td>3.2</td>
<td>1.2</td>
</tr>
<tr>
<td></td>
<td>2.0±0.1 [11]</td>
<td>7.2±0.5 [11]</td>
<td>2.7</td>
<td>1.0</td>
</tr>
<tr>
<td></td>
<td>1.6±0.17 [8]</td>
<td>–</td>
<td>3.3</td>
<td>–</td>
</tr>
<tr>
<td></td>
<td>1.83±0.18 [12]</td>
<td>6.77±0.25 [12]</td>
<td>2.9</td>
<td>1.1</td>
</tr>
<tr>
<td></td>
<td>3.0 [10]</td>
<td>7.3 [10]</td>
<td>2.5</td>
<td>1.3</td>
</tr>
<tr>
<td>$\nu_4$</td>
<td>1.5±0.1</td>
<td>4.9±0.9</td>
<td>3.5</td>
<td>1.6</td>
</tr>
<tr>
<td></td>
<td>1.5±0.1 [11]</td>
<td>6.4±0.7 [11]</td>
<td>3.5</td>
<td>1.1</td>
</tr>
<tr>
<td></td>
<td>1.2±0.17 [8]</td>
<td>–</td>
<td>4.4</td>
<td>–</td>
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<tr>
<td>AN solution</td>
<td></td>
<td></td>
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<tr>
<td>1-decyl-3-methyl-imidazolium bromide $^*$</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\nu_2$, 10%</td>
<td>3.4 (3.4)**</td>
<td>6.4 (–)**</td>
<td>1.6 (1.6)</td>
<td>1.8 (–)</td>
</tr>
<tr>
<td></td>
<td>4.4 (4.8)**</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td>20%</td>
<td>3.5 (2.7)</td>
<td>6.3 (6.7)</td>
<td>1.5 (2.0)</td>
</tr>
<tr>
<td></td>
<td>4.8 (4.9)**</td>
<td></td>
<td></td>
<td>1.9 (1.3)</td>
</tr>
<tr>
<td>$\nu_4$, 10%</td>
<td>2.7 (3.1)</td>
<td>– (–)</td>
<td>2.0 (1.7)</td>
<td>– (–)</td>
</tr>
<tr>
<td></td>
<td>20%</td>
<td>2.4 (2.6)</td>
<td>6.0 (6.1)</td>
<td>2.2 (2.0)</td>
</tr>
<tr>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.5 (1.5)</td>
</tr>
</tbody>
</table>

$^*$ measured at 295 and 329 K (in parentheses).

** the errors, after a nonlinear curve fitting procedure by Lorentz functions [9] was applied, and the extra bands were removed from the fundamentals (Figs. 3–5), do not exceed 10% for VV and about 20% for VH contours.

*** fit results using a single Lorentz contour given for comparison.

the present stage of work, i.e. before starting on AN dynamics in IL environment, it is only important to note that the half-widths and the values of $\tau_{vib}$ and $\tau_{rot}$ (2) obtained by digital processing of $\nu_2$ and $\nu_4$ band profiles pertain well to the data available in literature. They are collected in Table 1. Hence, the instrumental settings used in this study (see Sec. 2) are likely to be ‘metrologically’ approved.

Some of the experimental Raman contours of AN dissolved in 1-decyl-3-methyl-imidazolium bromide are shown in Figs. 4 and 5. The measured half-widths (HWHH) of Raman modes $\nu_2$ and $\nu_4$ and calculated vibrational and rotational correlation times ($\tau_{vib}$ and $\tau_{rot}$) at $T = 295$ and 329 K are given in Table 1. Unfortunately the dilution of 10 wt% of AN was too strong to extract the shapes of $\nu_4$ mode in VH polarization from the noise. The same happened with the $\nu_2$ contour at the higher temperature.

Secondary structures of the $\nu_2$ and $\nu_4$ Raman bands seen in the neat AN (Fig. 3) disappear in certain solutions, e.g. in water, formic acid, water/polymer gels [10, 13]. With increasing dissolution the overlap of bands increases, the sum-contour loses the asymmetry, and in the sequel the overlapping band components can not be resolved by statistical processing using least-squares fitting procedures [13]. However, more or less pronounced secondary structure of these bands is preserved in the investigated AN/IL solutions (Figs. 4 and 5). The VV spectra of $\nu_2$ and $\nu_4$ modes were fitted to the sum of two Lorentz profiles using the nonlinear regression method based on the Levenberg–Marquardt algorithm [9]. An enrolment of the second Lorentz function into a fitting procedure led to significant improvements of the statistical parameters $\chi^2$ (the square deviation) and $R^2$ (the correlation coefficient). Some of those cases are demonstrated in Figs. 4 and 5. The convergence to the same values of contour parameters (maximum position, HWHH, etc) independently of their initial values (zero-order approach) was achieved practically in all cases. Unstable fit was met for 10% solution at the lower temperature (i.e. 295 K) only. The overlap of VV contours at these conditions indeed seems to be too strong to allow to resolve them. The band fitting with a single Lorentz function gave overestimated HWHH value (Table 1) that differed from the true widths sufficiently. Therefore an approximate value of $\Delta_{VV} = 3.4$ cm$^{-1}$, i.e. the same as in 20% solution (Fig. 4), was used in the calculation of $\tau_{vib}$ and
Because of noise all investigated VH contours of AN solutions in IL were processed using a single Lorentz function approach. A certain systematic error estimating HWHH is very likely. Nevertheless this error contributes to all estimations of band parameters of VH contours, and thus qualitative changes in the rotational dynamics of AN molecules can be deduced.

The main conclusion is that the vibrational relaxation processes are quickened in the IL surrounding. The corresponding correlation times ($\tau_{\text{vib}}$, Table 1) decrease from about 3 ps in the neat AN to 1.5–2 ps in AN/IL solutions. The origin of this phenomenon is not yet clear. The IR studies on the isotopically diluted AN (5 mol % AN-$h_3$/AN-$d_3$, [11]) suggested that the contribution of the resonance energy transfer is rather small for the AN system. Its contribution to $\nu/2$ was not observed at all. Hence, some new mechanisms, specific to ILs due to their local ionic structuring, which gain the vibrational relaxation processes, may be expected. One of those mechanisms can be the intermode exchange of vibrational energy. This sort of dephasing does not lead to the dissipation of energy from the vibrational subsystem, but only causes its redistribution among additional vibrational degrees of freedom. Quite strong intermolecular interactions of the probe molecule with polyatomic ionic species of IL (in the present case – AN and the cationic part of 1-decyl-3-methyl-imidazolium bromide, respectively) can certainly favour this process.

The effect of IL surrounding on the rotational dynamics of AN, namely the tumbling motion of dipole moments is rather small. Although a certain slow-down of AN reorientations in IL can be distinguished, i.e. $\tau_{\text{rot}}$ increases from 1.0–1.5 ps (neat AN) to 1.6–1.9 ps (AN in IL), yet this effect is much less significant than it could be expected for so highly viscous media. For comparison, $\tau_{\text{rot}}$ values up to 2–4 ps were found for AN dissolved in water and water/polymer gels [13]. Such
drastic slow-down in solvents, the molecules of which possess proton-donor properties, can be understood as a consequence of hydrogen bonding of AN with solvent molecules. In the present case of 1-decyl-3-methylimidazolium bromide it looks that this IL surrounding acts on the rotation of molecular probe as the system being not too ‘super-polar’. In other words, considering microscopic friction effects, investigated 1-decyl-3-methylimidazolium bromide falls in the range between water, which forms huge H-bond networks and still is more polar than most ILs, and AN, where molecular reorientations are hindered by dipole–dipole correlations within small clusters [14, 15].

4. Concluding remarks

First of all, note that our study on the vibrational and rotational relaxations of AN dissolved in 1-decyl-3-methylimidazolium bromide has shown that these processes are rather weakly influenced by IL environment. The vibrational relaxation is more significantly enhanced, which is noticeable from about 50% decrease of vibrational correlation times in IL solution. The tumbling motion of electric dipoles of AN is much less affected by the IL surrounding in comparison with AN reorientations in water and water/polymeric network systems. One more very important aspect of this study should be noted. The dilution of 10 wt % was too high for registration of some Raman VH bands, but on the other hand maybe it was too strong with respect to the expected structural peculiarities in IL systems. The demolition of structural effects (short-range order and mesoscopic structuring in the nano-scale) in 1-decyl-3-methylimidazolium bromide by the mutual penetration of AN and IL species in such a big amount of AN molecules could be suspected. The experimental work carried out in the present study opens the ways for a wider use of this Raman spectroscopy studying fundamental molecular processes in ionic liquids.

Acknowledgements

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References

Santrauka

Joniniai skysčiai (medžiagos, kurios sudarytos iš imtinai tik iš anijonų ir katijonų ir kurių lydymosi temperatūra žemesnė nei +100°C, toliau – JS) yra vienas naujausių ir sekmę plėtojamos dangufunkcinių medžiagų kūrimo pradžių. Savo unikalios savybės JS pademonstravo įvairiose aukštojojių technologijų srityse, iš kurių ypač paminėti: enzimų katalizė, proteinių sintezė, membranų technologijos, kuro celes, dažnai naudojamos Saules baterijos (DSSC), dujų jutikliai, nanostruktūravimas ir kt. Milžiniškame darbe ir publikacijų sraute aptinkamos kai kurios labai svarbios tyrimų sujungtys, pvz., 4 ir 2) formas, užregistruotas naudojant įvairias šviesos poliarizacijos kombinacijas. Virpesinės ir rotacinės relaksacijų koreliacijų trukmės apskaičiuojamos pagal formules: 

\[ \tau_{\text{vib}} = \frac{1}{2 \pi (\Delta_{\text{iso}})} \quad \text{ir} \quad \tau_{\text{rot}} = \frac{1}{2 \pi (\Delta_{\text{anis}} - \Delta_{\text{iso}})}, \]

čia \( \Delta_{\text{iso}} \) ir \( \Delta_{\text{anis}} \) yra atitinkamai izotropinių ir anizotropinių juostų pslapčių, kurie buvo įvertinti ištyrusį smailis formos. Nustatyta, kad AN molekulių virpesinė reorientacinė dinamika nuo 3 ps grynajame AN iki 1,5–2 ps AN / JS tirpale, esant 10 ir 20 sv. % AN koncentracijai. Rotacinė reorientacinė suteikė įvairiuose režimiuose: nuo 1,0–1,5 ps (AN) iki 1,6–1,9 ps (AN / JS). Gautieji duomenys aptarianti tarpmodinės energijos perkėlimą, klimpos, tarpmolekulinės us gyvavimo, struktūrinių bei dieliktrinių JS ypatumų mechanizmų.