FABRICATION AND ELECTRICAL PROPERTIES OF $Li_3Sc_{2-x}B_x(PO_4)_3$ SUPERIONIC MATERIALS *

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The Li₃Sc_{2-x}B_x(PO₄)₃ compounds, with x = 0-2, were synthesized by a solid state reaction and studied by X-rays. The ceramic samples were sintered and studied by digital scanning microscope. The electrical properties of the ceramic (x = 0, 0.5, 1) and glass (x = 2) samples were investigated in a frequency range $10^6-1.2 \cdot 10^9$ Hz and in a temperature range 300–600 K. Two regions of relaxational dispersion were found in conductivity spectra of the ceramics. They are related to the fast Li⁺ ion transport in the bulk and in the grain boundaries of the ceramics. Isomorphic substitution Sc \rightarrow B leads to a change of temperature of the phase transition to γ -phase (T_{γ}). The increase of stoichiometric factor x caused the increases of intragrain ($\sigma_{\rm g}$) and total ($\sigma_{\rm tot}$) conductivities, and also the decrease of their activation energies in the temperature range below T_{γ} .

Keywords: solid electrolyte ceramics, superionic ceramics, fast ion transport, ionic conductivity

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1. Introduction

It is known that two structural phase transitions take place in $Li_3Sc_2(PO_4)_3$ crystals in temperature range from 293 to 573 K [1, 2]. The phase transitions are caused by order-disorder processes in the lithium sublattice. The $Li_3Sc_2(PO_4)_3$ compound at room temperature (α -phase) and at T = 473 K (β -phase) belongs to a monoclinic crystallographic system (s. g. $P2_1/n$), however, at T = 573 K (γ -phase) the crystals belong to the rhombic spatial symmetry group Pcan [2, 3]. In the γ -phase the lithium sublattice is disordered the most and the γ -Li₃Sc₂(PO₄)₃ crystals are solid electrolytes with fast Li⁺ ion transport [2, 4]. The conductivity σ of $Li_3Sc_2(PO_4)_3$ crystals was investigated in the frequency range from 5 Hz to $5 \cdot 10^5$ Hz [1, 2, 4]. The highest value of ionic conductivity of γ -Li₃Sc₂(PO₄)₃ single crystals was observed along the \vec{c} axis and at T = 600 K was found to be $\sigma = 3.0$ S/m (activation

energy $\Delta E = 0.38$ eV) [2]. The temperature of the phase transitions, the values of σ and ΔE are sen-

sitive to the isomorphic substitutions $Sc \rightarrow Fe$ in the

 $Li_3Fe_xSc_{2-x}(PO_4)_3$ quasibinary system. As it was

shown in Refs. [1,5] the lithium subsystem in the

quasibinary system remains disordered down to room

temperature for $0.2 \ge x \ge 0.6$. In temperature range from

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elongs 293 to 600 K the value of ionic conductivity of the sys- $2_1/n$, tem is higher than that of $Li_3Sc_2(Fe_2)(PO_4)_3$ crystals. Since Li^+ -ion conductors have appeared as attractive materials for applications in CO₂ gas sensors and solid electrolyte batteries, the detection of the influence of the isomorphic substitutions of Sc^{3+} by another threevalence ion on the values of σ and ΔE of $Li_3Sc_2(PO_4)_3$ compounds stimulates further research of this phenomena. Our intention was to investigate the influence of substitution of Sc^{3+} by a comparatively small B^{3+} cation. We report the technological conditions of synthesis of the new compound $Li_3B_xSc_{2-x}(PO_4)_3$ powders with x = 0-2, preparation of the glass samples, sintering of the ceramics, and results of our investi-

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gation of structure, surface, and electrical properties of the materials.

2. Experimental

Powdered compounds $\text{Li}_3\text{Sc}_{2-x}\text{B}_x(\text{PO}_4)_3$ at x = 0-2 have been synthesized by subjecting to a solid phase reaction the stoichiometric mixtures of Li_2CO_3 (purity 99.999%), Sc_2O_3 (purity 99.999%), extra pure NH₄H₂PO₄, and H₃BO₃. Each mixture was milled for 8 h in an agate planetary mill in ethyl alcohol, heated for 20 h at temperature T = 723 K, and milled for 12 h in the planetary mill. The synthesis continued with the sequence of heating for 2 h at T = 1273 K and milling in the planetary mill for 8 h. The fine powder was dried at T = 393 K for 24 h.

The powder was uniaxially cold pressed at 300 MPa. Ceramic pellets were sintered in air for 1 h. The sintering temperatures of Li₃Sc₂(PO₄)₃, Li₃Sc_{1.5}B_{0.5}(PO₄)₃, and Li₃ScB(PO₄)₃ ceramic samples were 1643, 1043, and 1023 K, respectively. The liquid phase of Li₃B₂(PO₄)₃ compound (the liquid phase temperature was found to be T = 1273 K) was air-quenched with a cooling rate of about 300 degree/min. The Li₃B₂(PO₄)₃ sample fabricated in this way was a single-phase glass. The structure parameters were analyzed at room temperature using diffraction of Cu $K_{\alpha 1}$ radiation from powder in the 2 Θ region from 10 to 80 degrees with a step scan of 1 degree/min.

The studies of surfaces of the ceramic samples were carried out by the scanning electron microscope (SEM) 960 DSM. The measurements of complex conductivity $\tilde{\sigma} = \sigma' + i\sigma''$, complex impedance $\tilde{Z} = Z' + iZ''$, and complex dielectric permittivity $\tilde{\varepsilon} = \varepsilon' + i\varepsilon''$ were performed by coaxial impedance spectrometer set-up in the frequency range from 10^6 to $1.2 \cdot 10^9$ Hz [6].

3. Results and discussion

The results of X-ray diffraction studies have shown that $\text{Li}_3\text{Sc}_{2-x}\text{B}_x(\text{PO}_4)_3$ at x = 0, 0.5, and 1 are mainly single-phase materials. A small amount of AlPO₄ was



Fig. 1. SEM image of $Li_3Sc_{1.5}B_{0.5}(PO_4)_3$ surface.

detected in samples as impurities from the material of crucible. These compounds belong to the monoclinic crystallographic system (s. g. $P2_1/n$) with z = 4 formula units in the unit cell. The results of X-ray investigations revealed changes in the interplanar distances with change of x in these compounds. The lattice parameters, unit cell volume, and density of the materials are presented in Table 1.

The density of ceramic samples was found to be from 86 to 92% of the theoretical density of these materials. The results of investigation by SEM showed the presence of numerous microcracks on the surface of ceramic specimens. The characteristic image of SEM of $Li_3Sc_{1.5}B_{0.5}(PO_4)_3$ ceramic surface is shown in Fig. 1. The grains of ceramic samples sintered for 1 h have the sizes in the range from 300 to 700 nm. The results of investigation of the element composition of studied ceramics are shown in Table 2.

The temperature dependence of electric conductivity was derived from the impedance spectra measured at different temperatures. The results of measurements of the frequency dependences of Z'' and σ' show that the impedance spectra consist of two overlapping dispersion regions. Both processes are related to the ion

Compounds	Latt	ice parameter	Unit cell	Density,	
	a	b	С	volume, Å ³	g/cm^3
$Li_3Sc_2(PO_4)_3$	8.853(2)	12.273(2)	8.802(2)	956.36	2.747
$Li_3Sc_{1.5}B_{0.5}(PO_4)_3$	8.858(3)	12.290(3)	8.810(3)	959.17	2.619
$Li_3ScB(PO_4)_3$	8.849(1)	12.281(3)	8.799(1)	956.22	2.510
$Li_3B_2(PO_4)_3$			glass		

Table 1. Cell parameters and densities of the $Li_3Sc_{2-x}B_x(PO_4)_3$ compounds at T = 300 K.

Compound	Elements, wt %					
	Р	0	В	Al	Li	Sc
$Li_3Sc_2(PO_4)_3$	23.49	48.53	_	_	5.26	22.73
$Li_3Sc_{1.5}B_{0.5}(PO_4)_3$	24.55	50.72	1.43	2.72	5.50	17.81
$Li_3ScB(PO_4)_3$	25.70	53.11	2.99	3.23	5.76	12.44
$Li_3B_2(PO_4)_3$	28.39	58.65	6.61	3.00	6.36	-

Table 2. The element compositions on the surface of $Li_3Sc_{2-x}B_x(PO_4)_3$ ceramics.

transport within grains and in grain boundaries of the ceramic samples. The dispersion processes are thermally activated and shift towards higher frequencies with the increase in temperature. This shows that both dispersions have typical relaxation character. The observed relaxation dispersions can be analyzed using the frequency dependences of ac impedance in the complex plane.

As an example, the complex plane impedance plots of $Li_3Sc_{1.5}B_{0.5}(PO_4)_3$ ceramic sample at different temperatures are shown in Fig. 2. The plot consists of two overlapping semicircles with centres below the Z' axis. The relaxation processes within grains are responsible for the high frequency arc, and the relaxation processes in the grain boundaries of the ceramic samples are responsible for the low frequency arc.

With the increase in temperature the relaxation dispersion frequency range related to the processes within grains of the ceramics gradually moves out from the investigated frequency range. The values of the intragrain conductivity σ_g and total conductivity σ_{tot} were obtained from the dependences Z''(Z') and $\sigma''(\sigma')$ at different temperatures. The temperature dependences of σ_g and σ_{tot} are shown in Figs. 3 and 4, respectively. The activation energies of σ_g and σ_{tot} were calculated according to Arrhenius equation

$$\sigma_{\rm g,tot} = \frac{\sigma_0}{T} \exp\left(\frac{\Delta E_{\rm g,tot}}{kT}\right),\tag{1}$$

where $k = 1.38 \cdot 10^{-23} \text{ J} \cdot \text{K}^{-1}$ is Boltzmann constant. The $\sigma_{\text{g}}(T)$ dependence of Li₃Sc₂(PO₄)₃ in Fig. 3 can be divided into four regions: below 450 K, from 450 to 470 K, from 470 to 530 K, and above 530 K. According to the structural analysis results for Li₃Sc₂(PO₄)₃ crystals [2], the anomalies of σ_{g} in temperature ranges 450 - 470 K and 470 - 530 K can be caused by the $\alpha \leftrightarrow \beta$ and $\beta \leftrightarrow \gamma$ phase transitions in the investigated compounds.

The results of investigation of $\sigma_{g,tot}(T)$ have revealed that the isomorphic substitution $Sc \rightarrow B$ leads to the changes of temperature T_{γ} of phase transitions to the γ -Li₃B_xSc_{2-x}(PO₄)₃ phase and also influences the values of σ_g and σ_{tot} as well as their activation ener-



Fig. 2. Impedance spectra of $Li_3Sc_{1.5}B_{0.5}(PO_4)_3$ ceramics at different temperatures.

gies. The increase of the stoichiometric factor x caused the increase of σ_g and decrease of σ_{tot} , ΔE_g , ΔE_{tot} in the low temperature region, while in the γ -phase it caused the decrease of both σ_g and σ_{tot} . The value of σ_{tot} of Li₃B₂(PO₄)₃ glass samples at low temperature is higher than the value of σ_g of α -Li₃Sc₂(PO₄)₃



Fig. 3. Temperature dependences of intragrain conductivity of $Li_3Sc_{2-x}B_x(PO_4)_3$ ceramics at x = 0, 0.5, and 1.

compound. The temperatures of γ -phase transitions decrease with the increase of x. The temperatures of γ -phase transitions for the compounds with stoichiometric factor x = 0, 0.5, and 1 are around 530, 520, and 510 K, respectively. The anomalies of σ_g at $\alpha \leftrightarrow \beta$ phase transition were not detected in the compounds with stoichiometric factor x = 0.5 and 1. The values of $\sigma_g, \sigma_{tot}, \Delta E_g$, and ΔE_{tot} of the investigated compounds are listed in Table 3. The relaxation frequency of the intragrain processes in ceramic samples was obtained from the maximum of frequency dependence of Z'' at different temperatures. The characteristic frequency dependences of Z'' for Li₃Sc₂(PO₄)₃ and the inverse temperature dependences of the σ_g relaxation frequencies are shown in Fig. 5. The relaxation fre-



Fig. 4. Temperature dependences of total conductivity of $Li_3Sc_{2-x}B_x(PO_4)_3$ ceramic at x = 0, 0.5, and 1 and of $Li_3B_2(PO_4)$ glass samples.

quencies within the grains of ceramic samples increase with temperature according to the formula

$$f_{\rm g} = f_0 \exp\left(\frac{\Delta E_f}{kT}\right),$$
 (2)

where f_0 is an attempt frequency related to the lattice vibrations and E_f is the activation energy of f_g . The comparison of the estimated values of ΔE_g and ΔE_f shows that these energies are approximately equal for the same compound. Since we have found that the activation energies of intragrain ionic conductivity of the investigated ceramic samples are equal to the activation energies of relaxation frequency, which can be attributed to the migration of Li⁺, the concentration of charge carriers remains constant with temperature.



Fig. 5. (a) Frequency dependences of Z'' of γ -Li₃Sc₂(PO₄) compound at different temperatures. (b) Temperature dependences of the relaxation frequency of the dispersion σ_g of Li₃Sc_{2-x}B_x(PO₄)₃ ceramics at x = 0, 0.5, and 1.

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<i>Т</i> , К	Stoichiometric factor x	$\sigma_{\rm g},$ S/m	$\Delta E_{\rm g}, {\rm eV}$	$\sigma_{ m tot}, { m S/m}$	$\Delta E_{\mathrm{tot}}, \mathrm{eV}$		
400	0	$1.99 \cdot 10^{-4}$	0.76	- 2	_		
	0.5	$6.5 \cdot 10^{-3}$	0.57	$1.2 \cdot 10^{-3}$	0.64		
	1	$3.93 \cdot 10^{-3}$	0.53	$1.18 \cdot 10^{-3}$	0.59		
	2 (glass)	-	_	$7.7 \cdot 10^{-4}$	0.53		
460	0	$4.3 \cdot 10^{-3}$	1.32	_	-		
	0.5	0.064	0.57	$11.5 \cdot 10^{-2}$	0.64		
	1	0.36	0.53	$7.68 \cdot 10^{-3}$	0.59		
	2	-	_	$5.5 \cdot 10^{-3}$	0.53		
540	0	1.15	0.37	0.25	0.47		
	0.5	0.545	_	0.15	_		
	1	0.27	_	0.085	_		
	2	-	-	-	_		

Table 3. Values of σ_{g} , σ_{tot} , ΔE_{g} , ΔE_{tot} of Li₃Sc_{2-x}B_x(PO₄)₃ ceramics at different temperatures.



Fig. 6. Temperature dependences of dielectric permittivity of $Li_3Sc_xB_{2-x}(PO_4)$ ceramic at $x = 0, 0.5, and 1 and of Li_3B_2(PO_4)$ glass samples.

Such ion transport peculiarities are characteristic for oxygen vacancy [7, 8], Na⁺ [9], and Li⁺ [10, 11] solid electrolytes.

Temperature dependences of the real part of complex dielectric permittivity ε' of $\text{Li}_3\text{Sc}_{2-x}\text{B}_x(\text{PO}_4)_3$ compounds are shown in Fig. 6. The dielectric permittivity was obtained at the electric field frequency of 1 GHz. The values of ε' at 300 K were found to be $\varepsilon' = 7$ for $\text{Li}_3\text{Sc}_2(\text{PO}_4)_3$ and $\text{Li}_3\text{ScB}(\text{PO}_4)_3$ compounds, $\varepsilon' = 8$ for $\text{Li}_3\text{Sc}_{1.5}\text{B}_{0.5}(\text{PO}_4)_3$ compound, and $\varepsilon' = 11$ for $\text{Li}_3\text{B}_2(\text{PO}_4)_3$ glass. The $\alpha \leftrightarrow \beta$ and $\beta \leftrightarrow \gamma$ phase transitions in $\text{Li}_3\text{Sc}_2(\text{PO}_4)_3$ compounds can be seen as anomalies of ε' . In the compounds with x =0.5 and 1 the anomalies of ε' were detected at 520 and 510 K temperatures, respectively. The activation energies of σ_g and σ_{tot} of these compounds also suffer the changes at exactly the same temperatures (Figs. 3, 4). In the γ -phase the value of ε' can be related to the contribution of polarization due to Li⁺ migration, vibrations of lattice, and electronic polarization.

4. Conclusions

The solid electrolyte $Li_3Sc_{2-x}B_x(PO_4)_3$ compounds at x = 0-2 have been synthesized by solid state reaction and studied by X-ray powder diffraction. The surface of the ceramic samples was investigated by SEM. The electrical properties of the ceramics and $Li_3B_2(PO_4)_3$ glass samples were carried out in the frequency range 10^{6} -1.2·10⁹ Hz and in the temperature range 300-600 K. Two regions of relaxation dispersion of the conductivity were found in the ceramics. The dispersions are caused by fast Li⁺ ion transport in the bulk and in the grain boundaries of the ceramics. Isomorphic substitution $Sc \rightarrow B$ leads to the changes of temperature of phase transition to the γ -Li₃Sc_{2-x}B_x(PO₄)₃ phase. The increase of stochiometric factor x causes the increase of $\sigma_{\rm g}$ values and the decrease of $\sigma_{\rm tot}$ and activation energies of the conductivities in the temperature region below T_{γ} , but the decrease of both $\sigma_{\rm g}$ and $\sigma_{\rm tot}$ in the γ -phase. The temperature dependences of the bulk conductivity and relaxation frequency in the grains of γ -phase give the same value of the activation energy. This suggests that fast Li⁺ ion transport in the grains of γ -Li₃Sc_{2-x}B_x(PO₄)₃ compounds may be described by the temperature dependent mobility, while the concentration of mobile ions remains constant with temperature. The increase of x leads to the increase of ε' at the frequency of external electrical field of 1 GHz. The

values of ε' at room temperature are related to the contribution of migration of Li⁺ ions, lattice vibrations, and electronic polarization.

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$Li_3Sc_{2-x}B_x(PO_4)_3$ SUPERJONIKŲ GAMYBA IR JŲ ELEKTRINĖS SAVYBĖS

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Santrauka

Pateiktos technologinės superjoninių Li₃Sc_{2-x}B_x(PO₄)₃ (čia x = 0 - 2) medžiagų sintezės sąlygos, jų rentgenostruktūrinių ir elementinės sudėties tyrimų rezultatai. Pagaminus Li₃Sc₂(PO₄)₃, Li₃Sc_{1,5}B_{0,5}(PO₄)₃ ir Li₃ScB(PO₄)₃ junginių keramikas bei Li₃B₂(PO₄)₃ stiklą, jų paviršiaus mikrosandara buvo tirta skenuojančiu elekroniniu mikroskopu. Elektrinės junginių savybės buvo tirtos 10⁶ – 1,2·10⁹ Hz elektrinio lauko bei 300 – 600 K temperatūros intervaluose. Li₃Sc₂(PO₄)₃ keramikose 450 – 480 K ir 500 – 535 K temperatūros intervaluose vyksta atitinkamai $\alpha \leftrightarrow \beta$ ir $\beta \leftrightarrow \gamma$ faziniai virsmai, kurių metu reiškiasi laidumo ir dielektrinės skvarbos anomalijos. Nustatyta, kad stechiometrijos faktoriaus x didėjimas žemina fazinio virsmo į junginių superjoninę γ faze

temperatūrą T_{γ} . Tirtose keramikose aptiktos dvi relaksacinio tipo dispersijos, susijusios su Li⁺ jonų pernaša kristalituose ir tarpkristalitinėse terpėse. Be to, didėjant x, esant žemesnei už T_{γ} temperatūrai, didėja keramikų kristalitinis (σ_g) ir bendrasis (σ_{tot}) laidumai, o jų aktyvacijos energijos mažėja. Parodyta, kad šiose medžiagose judriųjų Li⁺ jonų tankis kristalituose nepriklauso nuo temperatūros, o kristalitinio laidumo temperatūrinį kitimą lemia jonų judrio temperatūrinė kaita. Pateiktos tirtųjų keramikų ir Li₃B₂(PO₄)₃ stiklo santykinės dielektrinės skvarbos ε' , išmatuotos 1 GHz dažnio elektriniame lauke, temperatūrinės priklausomybės. Tų keramikų kristalitų dielektrinės poliarizacijų indėliai, tiek ir stechiomerijos faktoriaus x didumas.